

EXCEPTIONAL PARAFERMIONS IN A HILBERT SPACE OVER AN ASSOCIATIVE ALGEBRA

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Every vector of a Hilbert space admitting multiplication by octonions generates a linear manifold over the reals of at most 128 dimensions, and the basis for this manifold can provide a faithful representation of C_7 , the Clifford algebra of order seven⁽¹⁾. Multiplication defined as an equivalence relation in a (minimal one-sided ideal recovers the structure of the non-associative octonion algebra, and the basic properties of Hilbert spaces with octonion multipliers can therefore be studied in the framework of Hilbert spaces over associative algebras⁽²⁾. The set of algebraically closed linear manifolds of Hilbert spaces of this type has a structure consistent with that of the calculus of propositions characterizing quantum mechanics, and is isomorphic to a Hilbert space over a field in which there are superselection rules⁽³⁾.

Let e_1, e_2, \dots, e_7 be the generating elements of the associative real Clifford algebra C_7 , satisfying $\{e_a, e_b\} = -\delta_{ab}$, $e_a^* = -e_a$, $(e_a e_b)^* = e_b e_a$. A minimal projection in the algebra is

$$P_0^+ = \frac{1}{16} (1 - e_1 e_2 e_3)(1 - e_5 e_1 e_6)(1 - e_6 e_2 e_4)(1 + e_1 e_2 e_3 e_4 e_5 e_6 e_7), \quad (1)$$

where the $+$ superscript refers to the last sign. A complete basis is provided by $\rho_{ij}^\pm = e_i P_0^\pm e_j^*$ ($i = 0, \dots, 7$, $e_0 = 1$), which satisfies $\rho_{ij}^\sigma \rho_{kl}^{\sigma'} = \delta_{\sigma\sigma'} \delta_{jk} \rho_{il}^\sigma$ ($e_1 e_2 \dots e_7$ is in the center). The equivalence relation $P_0^+ ab = P_0^+(ab)_0^+$, where $(a)_0^+ = \sum e_i a_i$, a_i real, defines the multiplication law for octonions in the form used by Gürsey and Günaydin⁽⁴⁾.

The right linear scalar product $(f, g) = (g, f)^*$ has values in C_7 , $(f, f) \geq 0$ and $\|f\|^2 = \text{tr}(f, f)$, where the trace is defined in a representation of C_7 . Two algebraically closed linear manifolds, M_1 and M_2 , are orthogonal if $\text{tr}((f_1, f_2)a) = 0$ for all $a \in C_7$, for every $f_1 \in M_1$, $f_2 \in M_2$. If this condition is weakened to include only $a \in \mathbb{C}(1, e_7)$, the Abelian subalgebra spanned over the reals by 1 and e_7 , one obtains linear manifolds closed over $\mathbb{C}(1, e_7)$ corresponding to the proposal of Günaydin and Gürsey⁽⁴⁾. In this case, it suffices that $\text{tr}(f, g) = \text{tr}((f, g)e_7) = 0$. Using the algebraic identity (we restrict ourselves henceforth to $H^+ = H \sum_i \rho_{ii}^+$) $f = \sum_{ij} f_{ij} e_{ij}$, these bilinear forms are, respectively,

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$$\sum_{k,i} (f_{ki}, g_{ki}) ; \sum \{ (f_{k0}, g_{k7}) - (f_{k7}, g_{k0}) + (f_{k1}, g_{k4}) - (f_{k4}, g_{k1}) + (f_{k2}, g_{k5}) - (f_{k5}, g_{k2}) + (f_{k3}, g_{k6}) - (f_{k6}, g_{k3}) \} . \quad (2)$$

For each value of the index k , these forms are invariant, respectively, under $S0(8)$ and $Sp(8)$. The symmetry of the full scalar product (for each k) is therefore $S0(8) \cap Sp(8) = U(4)$. Indeed, let us call

$$\begin{aligned} \psi_0^k &= f_{k0} - f_{k7} e_7 \\ \psi_\alpha^k &= - f_{k\alpha} - f_{k\alpha+3} e_7 \end{aligned} \quad (3)$$

so that

$$f = \sum_k e_k P_k (\psi_0^k + \sum_{\alpha=1}^3 \psi_\alpha^k e_\alpha) . \quad (4)$$

Then,

$$(f, g)_C = \text{tr}(f, g) + e_7 \text{tr}((f, g)e_7) = \sum_k \{ (\psi_0^k, \chi_0^k) + \sum_\alpha (\chi_\alpha^k, \psi_\alpha^k) \} , \quad (5)$$

where χ corresponds to g . For each k , (5) is the scalar product given by Günaydin⁽⁴⁾. Observables may be defined to be linear over $\mathbb{C}(1, e_7)$, i.e., $A(fa) = (Af)a$, $a \in \mathbb{C}(1, e_7)$. The most general operator linear over $\mathbb{C}(1, e_7)$ satisfies⁽⁵⁾

$$Af = \sum_{k,j} e_k P_k \{ A_{00}^{kj} \psi_0^j + A_{0\alpha}^{kj} \psi_\alpha^{j*} + \sum_\alpha (A_{\alpha 0}^{kj} \psi_0^{j*} + \sum_\beta A_{\alpha\beta}^{kj} \psi_\beta^j) e_\alpha \} , \quad (6)$$

explicitly exhibiting superselection rules when H is viewed as a family of Hilbert spaces over $\mathbb{C}(1, e_7)$ (note that A may act antilinearly on this family). With the help of the form of (6), we may construct the tensor product

$$\begin{aligned} gTf &= \sum_j e_j P_j \{ (\chi_0^j + \chi_\alpha^j e_\alpha) \times (T_{00} \psi_0^j + T_{0\beta} \psi_\beta^{j*}) \\ &+ \sum_\alpha (\chi_\alpha^j + \chi_0^j e_\alpha) \times (T_{\alpha 0}^* \psi_0^j + T_{\alpha\beta}^* \psi_\beta^{j*}) \} , \end{aligned} \quad (7)$$

in which T acts linearly with a numerical matrix, and g acts linearly with χ_α^j occurring through an action analogous to the second and third terms of (6), and the first and fourth carry χ_0^j . Clearly, $gT(fa) = (gTf)a$, $a \in \mathbb{C}(1, e_7)$. For $\hat{T}_{0\alpha} = \hat{T}_{\alpha 0} = 0$, $\hat{T}_{\alpha\beta} = \hat{T}_{00} \delta_{\alpha\beta} = e_7 \delta_{\alpha\beta}$, (7) takes on the form of the tensor product given by Günaydin⁽⁴⁾. In this case, there exists an operator A_{12} such that $A_{12}(g\hat{T}f) = (A_1 g)\hat{T}(A_2 f)$, provided that A_1 and A_2 do not connect to ψ_α^{j*} , ψ_0^{j*} .

When these conjugate components are involved, the existence of A_{12} implies that $T_{\alpha 0} = 0 \leftrightarrow T_{\alpha\beta} = 0$ and $T_{0\alpha} = 0 \leftrightarrow T_{00} = 0$. If we wish to impose the condition $(ga) \check{f} = (g\check{f})a$, then we must take $\check{T}_{\alpha 0} = \check{T}_{\alpha\beta} = 0$, so that only the first term of the right hand side of (7) occurs in $g\check{f}$.

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REFERENCES

1. H.H. Goldstine and L.P. Horwitz, Proc. Nat. Aca. Sci. 48, 1134 (1962); Math. Ann. 154, 1 (1964). For recent interest, see F. Gürsey, Orbis Scientiae Conference at Coral Gables, Jan. 1976, and references therein. See also, A. Barducci, F. Buccella, R. Casalbuoni, L. Lusanna and E. Sorace, Firenze preprint* Jan. 1977, where the isomorphism between the Jordan products of the Fermion operators of C_6 and of the six transverse split octonions is noted. The authors were not interested in the non-observability aspects (presumably related to non-associativity), and therefore the associative realization was chosen and the automorphism groups of some of the Clifford algebras were studied. We emphasize in what follows that the non-associative aspects are implicit in the Clifford algebras, and hence the study carried out by Barducci et al. is actually of greater generality.
2. H.H. Goldstine and L.P. Horwitz, Math. Ann. 164, 291 (1966).
3. L.P. Horwitz and L.C. Biedenharn, Helv. Phys. Acta 38, 385 (1965).
4. M. Günaydin and F. Gürsey, Jour. Math. Phys. 14, 1651 (1973); Phys. Rev. D9, 3387 (1974); M. Günaydin, Jour. Math. Phys. 17, 1875 (1976).
5. See, for example, the representations of automorphism groups given by C. Saclioglu, Phys. Rev., in press.

* Phys. Lett. 67B, 344 (1977).